

Lecture 5. Gravity field –modelling - sphere.

We use the development for T and integrate over sphere:

$$\begin{aligned}
 T(\varphi, \lambda, r) &= \sum_{i=2}^{\infty} \frac{GM}{r} \left(\frac{R}{r}\right)^i \sum_{j=-i}^i Y_{ij}(\varphi, \lambda) \frac{i-1}{i-1} \left(\frac{R^2}{R^2}\right) \Delta \bar{C}_{ij} \\
 &= \sum_{i=2}^{\infty} \frac{1}{r} \left(\frac{R}{r}\right)^i \sum_{j=-i}^i Y_{ij}(\varphi, \lambda) \frac{R^2}{i-1} \frac{1}{4\pi} \iint \Delta g(\varphi', \lambda', R) Y_{ij}(\varphi', \lambda') dS \\
 &= \frac{R}{4\pi} \iint \left[\Delta g(\varphi', \lambda', R) \sum_{i=2}^{\infty} \frac{1}{i-1} \left(\frac{R}{r}\right)^{i+1} \sum_{j=-i}^i Y_{ij}(\varphi, \lambda) Y_{ij}(\varphi', \lambda') \right] dS
 \end{aligned}$$

Stokes formula

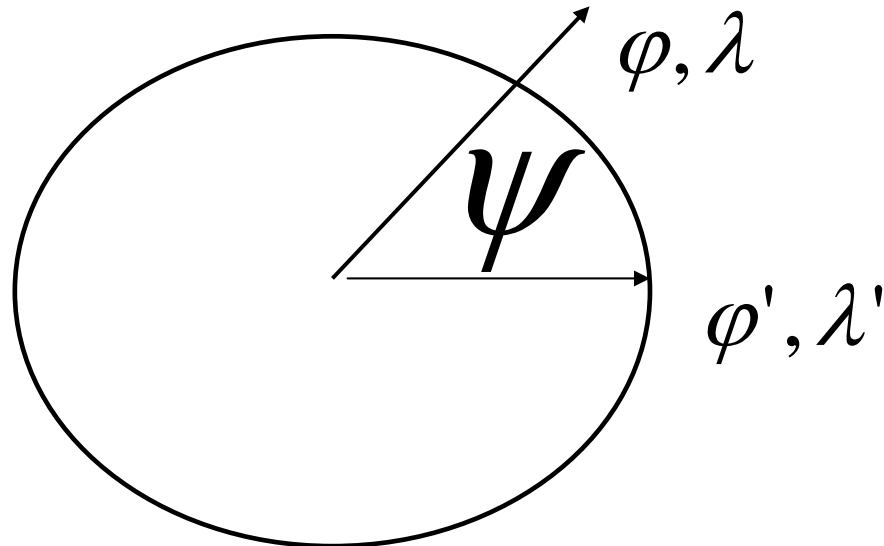
We get:

$$T(\varphi, \lambda, r) = \frac{R}{4\pi} \iint_S \Delta g \cdot \left(\frac{R}{r}\right)^{i+1} \sum_{i=2}^{\infty} \frac{2i+1}{i-1} P_i(\cos \psi) d\sigma$$

Gravity field-modelling.

We use the development in Legendre-polynomials

$$P_i(\cos \psi) = \frac{1}{2i+1} \sum_{j=-i}^i Y_{ij}(\varphi, \lambda) Y_{ij}(\varphi', \lambda')$$



Stokes formula, 1849, Pizetti, 1911.

We may obtain closed expression because:

$$\frac{1}{l} = \sum_{i=0}^{\infty} \frac{1}{r} \left(\frac{R}{r} \right)^i P_i(\cos \psi)$$

$$S(\psi, r) = \left(\frac{R}{r} \right)^{i+1} \frac{2i+1}{i-1} P_i(\cos \psi),$$

so on the sphere : $R / r = 1$

$$S(\psi, R) = \frac{1}{\sin(\psi / 2)} + 15 \cos \psi - 6 \sin(\psi / 2) \\ - 3 \cos \psi \cdot \ln(\sin(\psi / 2) + \sin^2(\psi / 2))$$

Stokes formula.

We have solved a boundary value problem for an elliptic partial differential equation, but

Note singularity:

$$\frac{1}{\sin(\psi / 2)} \rightarrow \infty$$

when

$$\psi \rightarrow 0$$

Integration must be done with special method.

Use of Stokes formulae.

Height-anomaly/geoid height:

$$\zeta(\varphi, \lambda, r) = \frac{R}{\gamma \cdot 4\pi} \iint_S \Delta g \cdot S(\psi, r) dS$$

Deflections of the vertical, Vening-Meinesz, 1928:

$$\xi(\varphi, \lambda, r) = \frac{1}{\gamma \cdot 4\pi} \iint_S \Delta g \cdot \frac{\partial}{\partial \varphi} S(\psi, r) dS$$

$$\eta(\varphi, \lambda, r) = \frac{1}{\gamma \cdot 4\pi \cos \varphi} \iint_S \Delta g \cdot \frac{\partial}{\partial \lambda} S(\psi, r) dS$$

Stokes formula.

Due to the singularity most of the contribution comes from data with a short distance from the computational-point.

A subtraction of a spherical harmonic model (EGM) to N=360, then only very small area is needed.

$$\zeta(\varphi, \lambda, r) = \zeta_{EGM}(\varphi, \lambda, r) + \frac{R}{\gamma \cdot 4\pi} \iint_{S_0} (\Delta g - \Delta g_{EGM}) \cdot S(\psi, r) dS$$

Note: EGM must be added back later !

Stokes formula, plane Earth.

If we want to compute the integral locally, we may regard the Earth as plane.

Integral may be computed using Fourier-transformation (convolution) in the plane, where the spectral relation $(i-1)/r$ is used.

$$\zeta(x, y, R) = \zeta(z, R)_{EGM} + \int \Delta g(t) S(z - t) dt$$

Developed by R.Forsberg and M.Sideris in 1986.

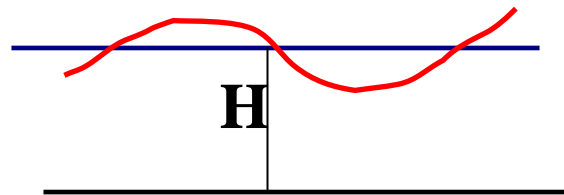
R.Forsberg/CCT: subtract (and add) also local topography

Stokes on the sphere.

Remove masses T_M over altitude 0 (add back later)

Gravity anomaly must then be computed in altitude zero 0
("downward continuation")

Simple method to remove: Bouguer-plate removed in each
point:

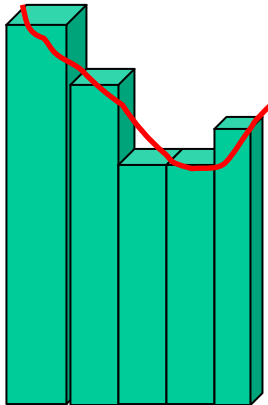


$$2\pi G\rho H \approx 0.11 \text{ mgal} / m$$

Not good because $T - T_M$ not harmonic anymore.

Removal of residual topography, Torge 6.5.

Rectangular prisms better:



May be refined with sloping boundaries – more details in R.Forsberg's lectures.

Boundary value problem for T .

For Harmonic functions, the functionen may be determined from its values on the boundary

Chauchy's boundary value problem

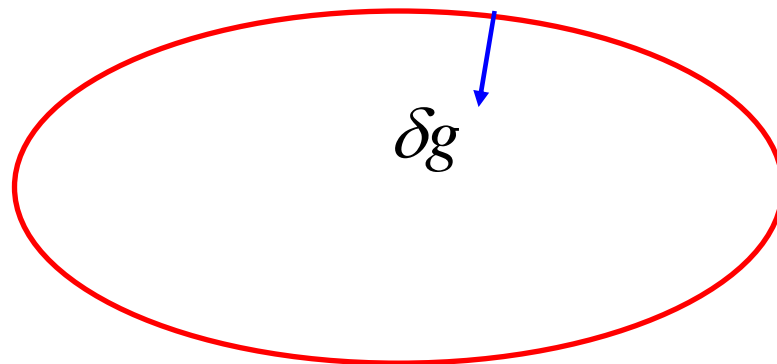
Or from the normal-derivative on the surface

Neumanns-randværdi-opgave

Side condition:

$\rightarrow 0$ for

$r \rightarrow \infty$



Pizettis formula.

We use:
$$\frac{GM}{R} \Delta \bar{C}_{ij} = \frac{1}{4\pi} \iint_S T(\varphi', \lambda', R) \cdot Y_{ij}(\varphi', \lambda') dS$$

$$\begin{aligned} T(\varphi, \lambda, r) &= \sum_{i=2}^{\infty} \frac{GM}{r} \left(\frac{R}{r}\right)^i \sum_{j=-i}^i Y_{ij}(\varphi, \lambda) \Delta \bar{C}_{ij} \\ &= \sum_{i=2}^{\infty} \frac{1}{r} \left(\frac{R}{r}\right)^i \sum_{j=-i}^i Y_{ij}(\varphi, \lambda) \frac{1}{4\pi} \iint T(\varphi', \lambda', R) Y_{ij}(\varphi', \lambda') dS \\ &= \frac{1}{4\pi} \iint \left[T(\varphi', \lambda', R) \sum_{i=2}^{\infty} \left(\frac{R}{r}\right)^{i+1} \sum_{j=-i}^i Y_{ij}(\varphi, \lambda) Y_{ij}(\varphi', \lambda') \right] dS \end{aligned}$$

Poissons Integral.

We use Legendre-polynomials expressed using Legendre-functions, and the summation formula for $1/(\text{distance})$:

$$T(\varphi, \lambda, r) = \frac{R^2 (r^2 - R^2)}{4\pi} \iint_S \frac{T(\varphi', \lambda', R)}{l^3} dS$$

Solves Cauchy's boundary value-problem for sphere for harmonic function, regular at infinity.

In the same way are $r \cdot \Delta g$ harmonic and may be upward continued or gives integral-equation so that gravity on the sphere can be determined.

Evaluation of the integrals.

(1) T or Δg known or “estimated” = “Predicted” everywhere on the sphere.

Collocation useful method !

(2) Possible to use functional connection between data and T for a direct estimation of T using collocation

$$T(P) = \sum_{i=1}^N a_i \cdot COV(T(P), observation(i)) =$$
$$\{observation(i)\}^T \{COV(obs(i), obs(j) + \sigma_{ij})\}^{-1}$$
$$\{COV(T(P), observation(i))\}$$

Solution in Hilbertspace.

Hilbertspace of harmonic functions has reproducing kernel (identity-mapping):

$$T(P) = (T(Q), K(P, Q)),$$

$(\dots, \dots) = \text{inner product in space}$

If F_i orthonormal basis, then

$$T(P) = \sum_{i=0}^{\infty} C_i \cdot F_i(P)$$

$$K(P, Q) = \sum_{i=0}^{\infty} F_i(P) \cdot F_i(Q)$$

Solution in Hilbertspace.

Equivalent with statistical Collocation there is a functional-analytic formulation, $\text{COV}(T(P), T(Q)) = K(P, Q)$.

We want to find an approximation, which agree with noise-free data:

$$L_i(T) = L_i(\tilde{T}) = \textit{observation}(i) \quad (\text{without noise})$$

Solution with minimum norm in the Hilbert-space:

$$\tilde{T}(P) = \sum_{i=1}^N b_i L_i(K(P, \cdot))$$

$$\{b_i\} = \{L_i L_j K(\cdot, \cdot)\}^{-1} \{\textit{observation}(j)\}$$

Isotropic inner product

R - rotation of the full space. Inner product isotropic if $(F,G)=(F(R),G(R))$. Reproducing kernel then like covariance function:

$$K(P, Q) = K(r, r', \psi) = \sum_{i=0}^{\infty} \sigma_i \left(\frac{R^2}{rr'} \right)^{i+1} P_i(\cos \psi)$$

Isotropic inner product, L_2 on the sphere.

Basis-functions: Normalized spherical harmonics.

$$K(P, Q) = K(\varphi, \lambda, r, \varphi', \lambda', r') =$$
$$\sum_{i=0}^{\infty} \left(\frac{R^2}{rr'} \right)^{i+1} \sum_{j=-i}^i Y_{ij}(\varphi, \lambda) Y_{ij}(\varphi', \lambda') =$$
$$\sum_{i=0}^{\infty} \left(\frac{R^2}{rr'} \right)^{i+1} (2i+1) P_i(\cos \psi)$$

Not usefull, because infinite on the sphere for distance zero.

= variancen of potentialet is infinite on the sphere.

Closed expression for kernel possible !

Degree-variances which give finite values of potential and gravity variance on sphere:

$$\sigma_i = \left(\frac{R_B^2}{R^2} \right)^{i+1} \frac{A}{(i-1)(i-2)(i+24)}$$

$R_B < R$, radius for Bjerhammar-sphere

Collocation with parameters.

Data may be biased – i.e. a constant or linear function has been added to data due to errors in data-collection or because local height datum has been used.

l = observation vector, X parameter vector

A = linearized observation – matrix

s = signal (gravity field),

$L(s)$ observation vektor, n = noise

$$l = AX + L(s) + n$$

Collocation with parameters.

Covariances or functionals applied on Kernel:

$$\bar{C} = \{COV(obs_i, obs_j) + \sigma_{ij}\}, n \times n \text{ matrix}$$

$$C_P = \{COV(T(P), obs_i)\}, n \text{ vector}$$

$$\tilde{X} = (A^T \bar{C}^{-1} A)^{-1} (A^T \bar{C}^{-1} l)$$

$$\tilde{s}(P) = \tilde{T}(P) = C_P^T \cdot \bar{C}^{-1} \cdot (l - AX)$$

Error-estimates may also be computed !

Integrated geodesy.

**One may mix gravity data with
geometrical data.**

Geodetic Earth-models

If only finite number of coefficients used (N) for T and M other data,

then we get traditional least-squares solution with $N+M$ unknowns. Typisk $N=73*73$ eller $361*361$.

$M=300*3+100$ tide-parametre+GM, a , f +datum-shift.

WGS84: NIMA (US DoD),

GEM: Goddard (NASA) Earth Model serie

GRIM: French/German model series

EIGEN: German/French/US

Only gravity: OSU91, EGM96, EGM2008.